

QCD and Hadronic interactions

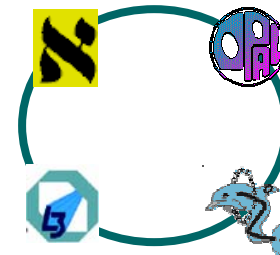
Recontres de Moriond-La Thuille 12-19 March 2005

Searches for Higgs in the MSSM CP-conserving and CP-violating scenarios at LEP

P. Ferrari (CERN)

on behalf of the LEP experiments

- model introduction
- the CP-conserving MSSM
- the CP-violating MSSM
- 2HDM



2 Higgs Doublet Models

- Simplest extension of SM are 2HDMs

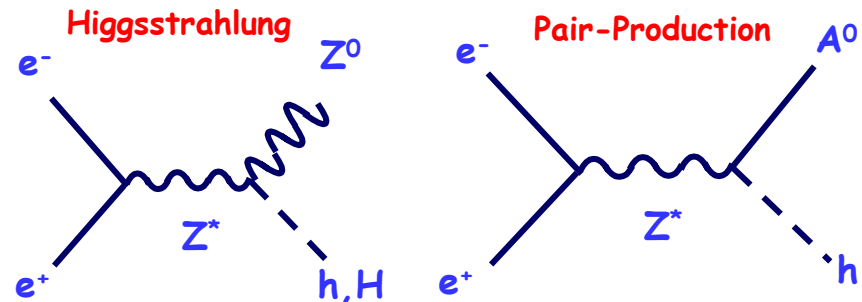
$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_w} \sim 1 \quad \text{and no FCNC}$$

- Two complex scalar field doublets Φ_1 and Φ_2 , 5 scalar Higgses:
 - Real parts mix with $\alpha \rightarrow$ CP even scalars h^0, H^0
 - Imaginary part \rightarrow CP odd scalar A^0
 - Two charged scalars H^{\pm}

- Two production processes :

- $\sigma_{hZ} = \sin^2(\beta - \alpha) \sigma_{HZ}^{SM}$
- $\sigma_{HZ} = \cos^2(\beta - \alpha) \sigma_{HZ}^{SM}$
- $\sigma_{HA} = \cos^2(\beta - \alpha) \lambda \sigma_{HZ}^{SM}$

β = ratio of VEV of scalar fields



- The type of 2HDM determined by the couplings of Φ_1 , Φ_2 to fermions:
 - Only Φ_1 couples to fermions 2HDM (I)
 - Φ_1 (Φ_2) couples to down (up) type fermions 2HDM(II)

The MSSM

- 2HDM(II) are interesting since by adding supersymmetry



CP-conserving MSSM

- CP-conserving MSSM is interesting since it provides framework for unification of Gauge interactions and stability of universe at EW scale

$m_h < 140 \text{ GeV}$ after radiative corrections

- to explain matter-antimatter asymmetry in universe we need



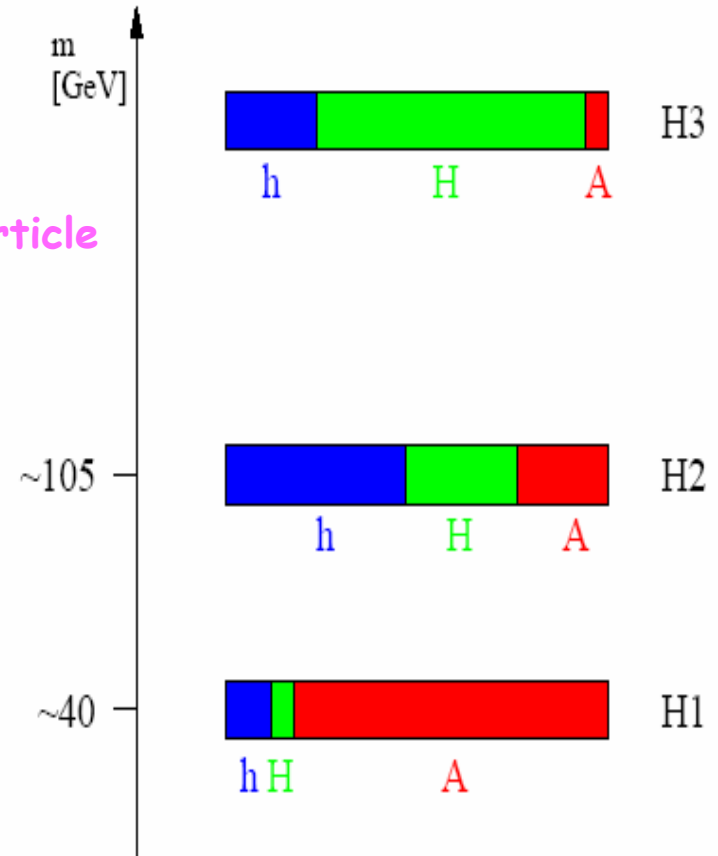
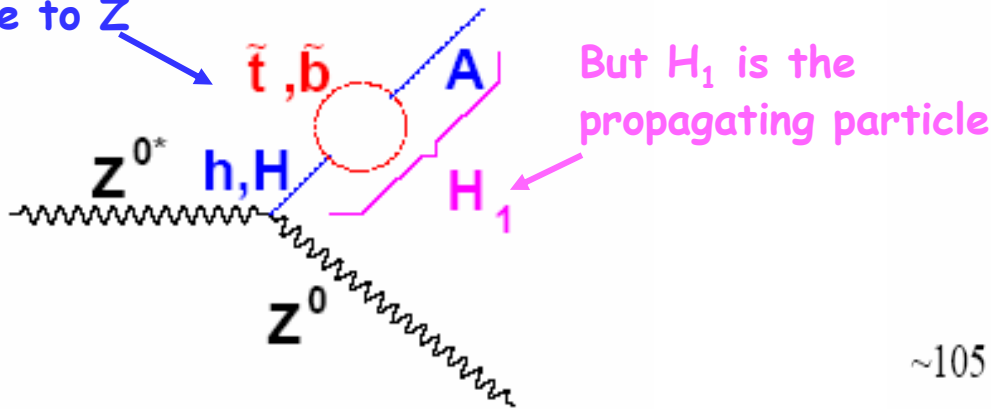
CP-violation \gg than in SM

Justifies introduction of CP-violation in MSSM: can be done via radiative corrections (in particular from 3rd generation s-quarks)

CP-violation in MSSM

Mass eigenstates and CP-eigenstates do not coincide:
 H_1, H_2, H_3 are mixtures of CP-even and CP-odd Higgs fields

Only CP eigenstates h, H
 can couple to Z



- Lightest Higgs boson might have escaped detection at LEP2: H_1 might decouple almost completely from the Z
- Search for both $H_1 Z$ and $H_2 Z$ production

Experimental searches

b-tagging HZ

- $(H \rightarrow bb) (Z \rightarrow qq)$
- $(H \rightarrow bb, \tau\tau) (Z \rightarrow \nu\nu)$
- $(H \rightarrow bb, qq) (Z \rightarrow ee, \mu\mu)$
- $(H \rightarrow \tau\tau) (Z \rightarrow qq), (H \rightarrow bb, \tau\tau) (Z \rightarrow \tau\tau)$

ALEPH, DELPHI, OPAL & L3
data @ $91 \text{ GeV} < \sqrt{s} < 209 \text{ GeV}$
Interpreted in MSSM
CPC and CPV and 2HDM(II)

Flavour independent HZ

- $(H \rightarrow qq) Z$
- $(H_2 \rightarrow H_1 H_1) Z$ dominant when kinematically allowed

b-tagging pair production $H_2 H_1$

- $(H_2 \rightarrow bb) (H_1 \rightarrow bb)$
- $(H_2 \rightarrow \tau\tau) (H_1 \rightarrow bb)$
- $(H_2 \rightarrow H_1 H_1 \rightarrow bbbb) ((H_1 \rightarrow bb))$

Flavour independent $H_2 H_1$ used only for 2HDM(II) scan

- $(H_2 \rightarrow qq) (H_1 \rightarrow qq)$

Additional constraints

- Z width, decay mode independent search for HZ, light Higgs from Yukawa production

MSSM CPC benchmarks

Traditional scans:

- **No mixing:** in stop sector
- **m_h max:** yields maximal bound on m_h^{TH}
- **Large μ :** suppressed $h \rightarrow bb$

7 parameters: (Carena et al. hep-ph/9912223)

$m_{\text{top}} = 179.3 \text{ GeV}$ ($178.0 \pm 4.3 \text{ GeV}$ CDF & D0)

M_{SUSY} sfermion mass at EW scale

μ Higgs mixing parameter

M_2 gaugino mass at EW scale

m_g gluino mass

X_t = Stop mixing parameter

$A_b = A_t = X_t + \mu \cot \beta$ = trilinear Higgs-squark coupling

New scans: envisaged for final LEP combination but not yet done

● Favoured by $(g-2)_\mu$ and $\text{Br}(b \rightarrow s\gamma)$

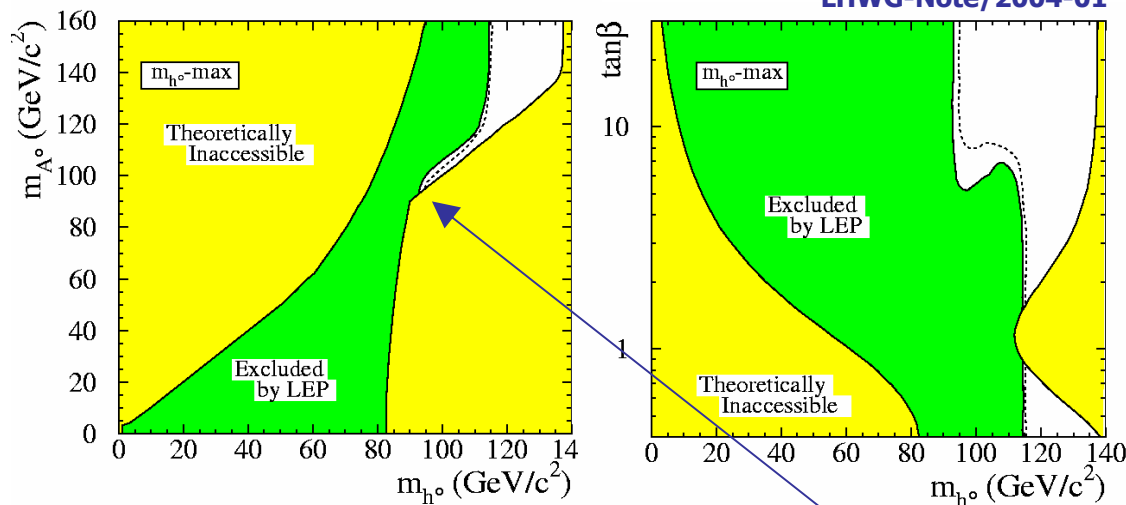
- **No mixing (2TeV)** reversed μ sign motivated by $(g-2)_\mu$
- **m_h -max+** with reversed μ sign motivated by $(g-2)_\mu$
- **constrained m_h -max** reversed sign for A_t and X_t motivated by $\text{Br}(b \rightarrow s\gamma)$

● Regions where Hadron colliders might have problems in detecting the Higgs

- **gluophobic** $gg \rightarrow h$ suppressed
- **small α_{eff}** $h \rightarrow bb, \tau\tau$ suppressed

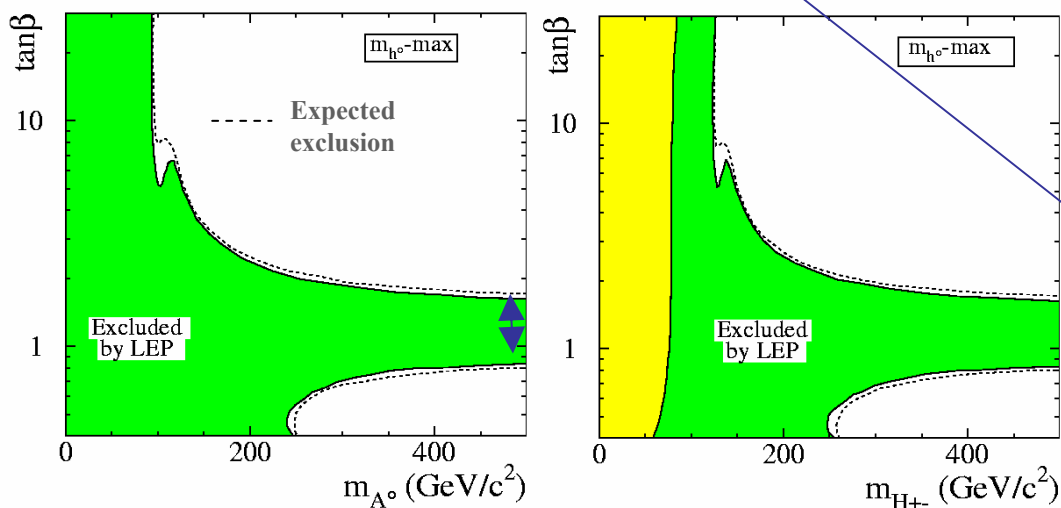
CPC m_h -max scan

LHWG-Note/2004-01



2 calculations used:

- FeynHiggs 2.0: 2-loop diagrammatic approach & OS scheme
S.Heinemeyer et al. hep-ph/0212037
- SUBHPOLE: 1-loop renormalization group, \overline{MS} scheme
M.Carena et al hep-ph/9912223



FeynHiggs is chosen:

more accurate, conservative results

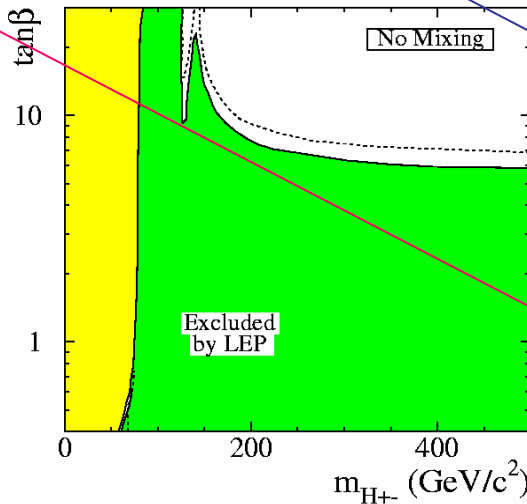
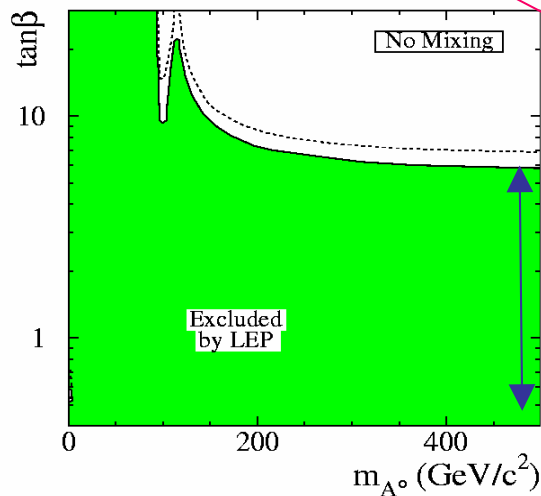
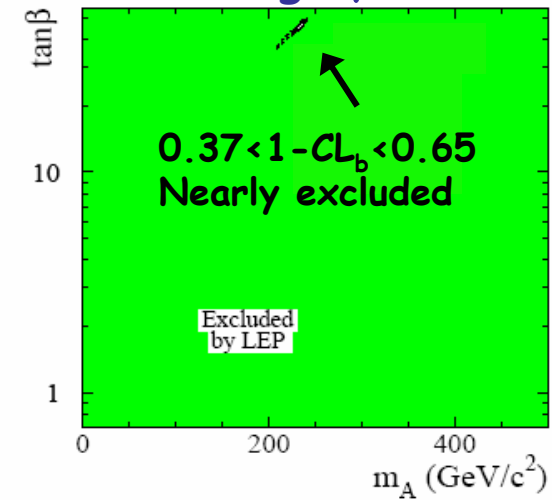
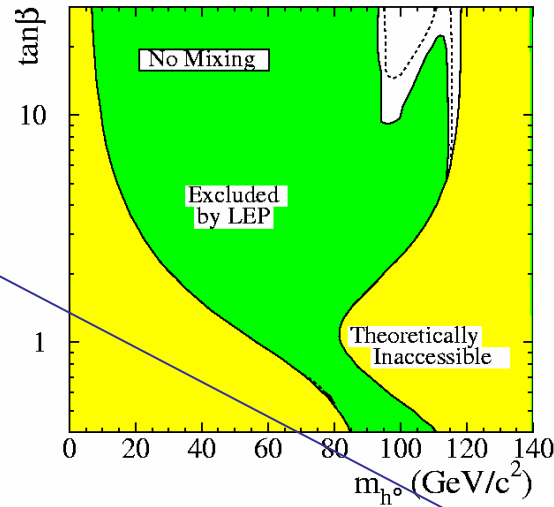
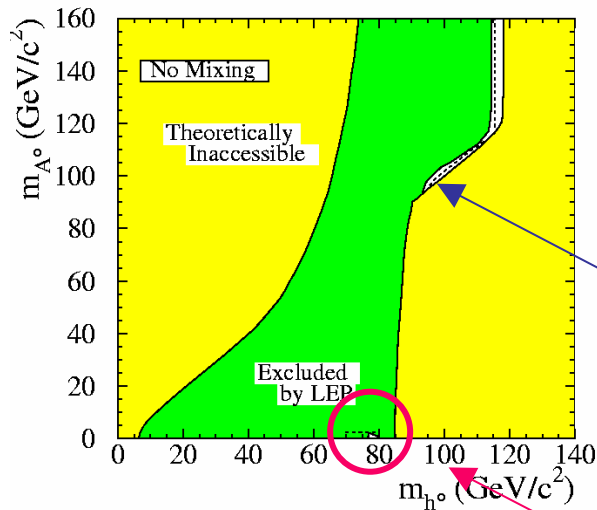
	m_h (GeV/c ²)	m_A (GeV/c ²)	$\tan\beta$
Obs	92.9	94.8	0.9-1.5
Exp	94.8	95.1	0.8-1.6

No-mixing & large μ scans

LHWG-Note/2004-01

No-mixing

Large μ

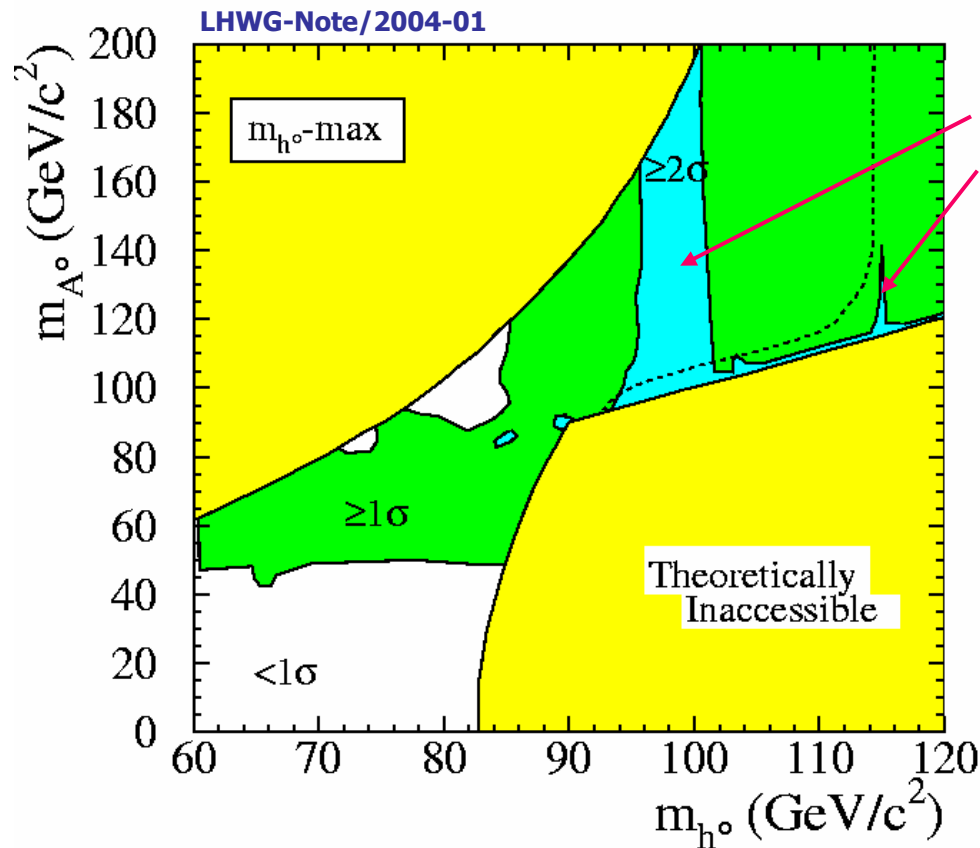


No-mixing

	m_h (GeV/c ²)	m_A (GeV/c ²)	$\tan\beta$
Obs	93.3	93.3	0.4-5.6
Exp	95.0	95.0	0.4-6.5

For $m_h \sim 80$ GeV and $\tan\beta < 0.7$, $m_A < \tau$ threshold:
A decays uncertain
 $h \rightarrow b\bar{b}$ small

Did we miss the Higgs?



No excess larger than 3σ :

- $\geq 2\sigma$ @ $m_h \sim 98$ GeV
- $\geq 2\sigma$ @ $m_h \sim 115$ GeV.

Recent interpretations:

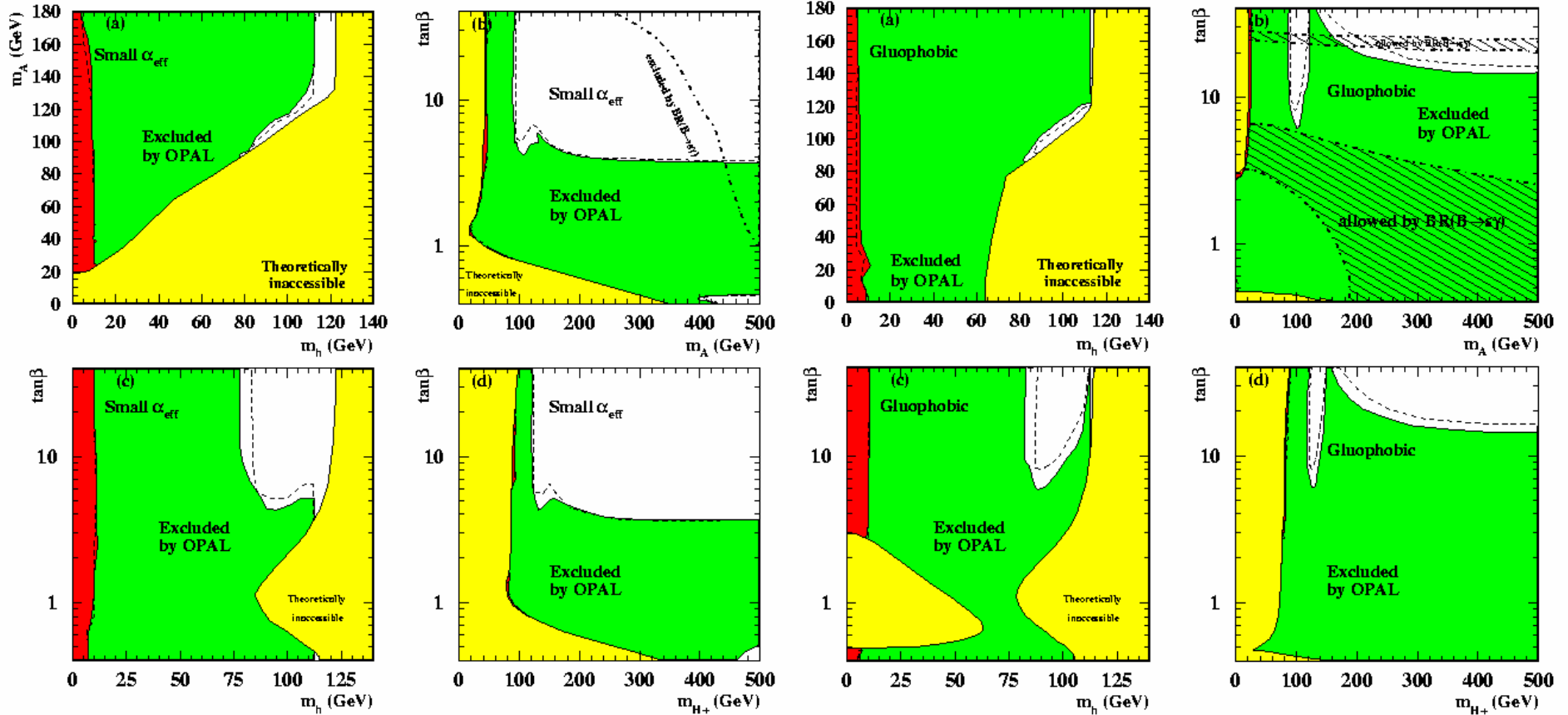
M. Drees hep-ph/0502075
G.L. Kane et al. hep-ph/0407001

Explain this kind of "excess" within

- CP-conserving MSSM
- CP-violating MSSM
- 2HDM

Example of new scans

OPAL Eur.Phys.J.C37:49-78,2004



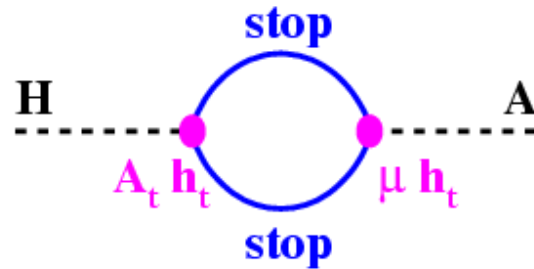
Only OPAL data: exclusion will be larger for LEP combination



Large regions of the parameters space are excluded

CP-violating MSSM

Phases of A_t , A_b and $m_{\tilde{g}}$ introduce CP violation in the Higgs potential via loop effects leading off-diagonal contributions to Higgs mass matrix



Theoretically: $\arg A_u \neq 0$ most general case, can be motivated by Baryogenesis.

Size of CP violating effects proportional to:

$$M_{SP}^2 \propto \frac{m_t^4}{v^2} \frac{\text{Im}(\mu A_t)}{32\pi^2 m_{SUSY}^2}$$

- benchmark: large $\arg A_u \neq 0$, large μ , relatively small m_{SUSY}
- the CP violation increases with m_t

CPX benchmark

Carena et al., Phys.Lett B495 155(2000)

$\tan\beta$	m_{H^\pm} (GeV)	μ (GeV)	m_{SUSY} (GeV)	M_2 (GeV)	$ A_q $ (TeV)	$\arg(A_q)$	m_g (TeV)	$\arg(m_g)$
0.6–40	4–1000	2000	500	200	1	90°	1	90°

EDM measurements of n and e fulfilled

Maximal CP violation

Feynhiggs and CPH are a priori equivalent:

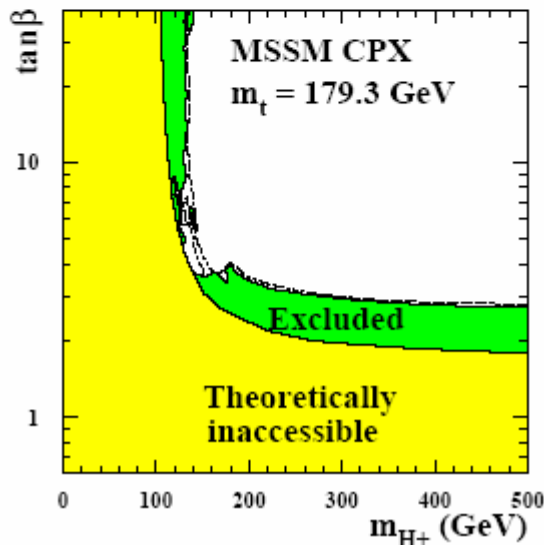
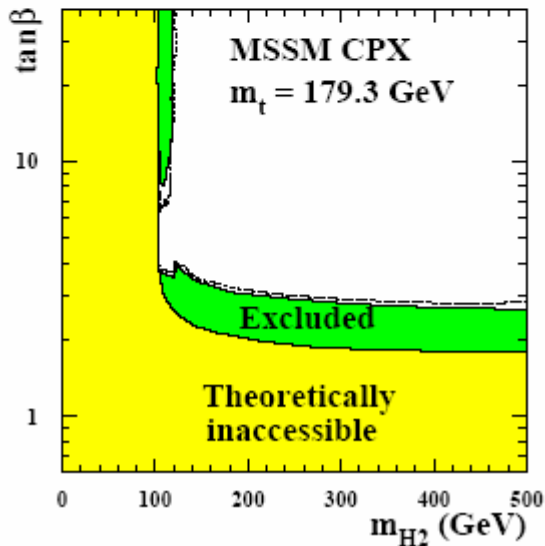
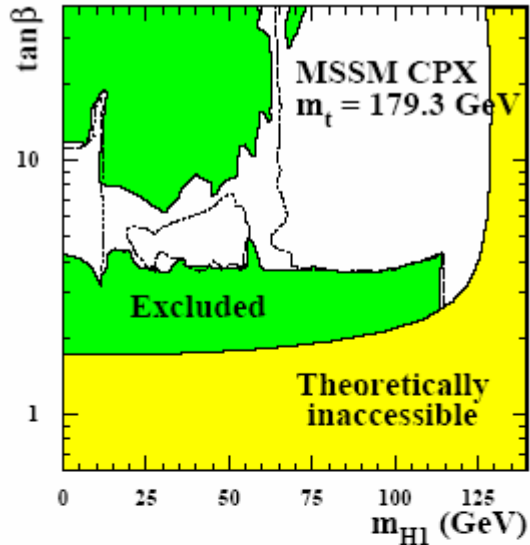
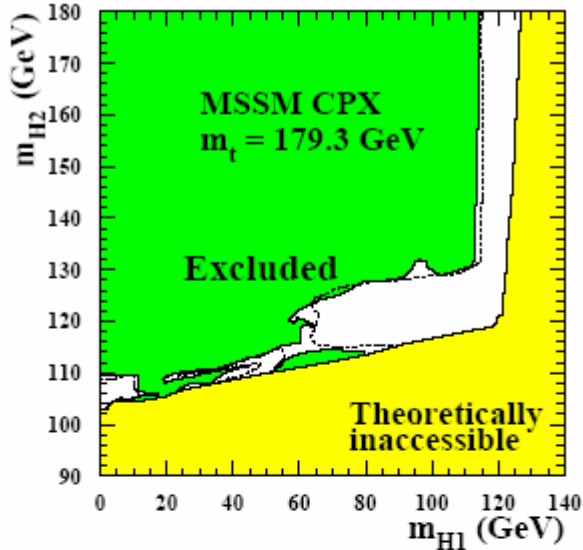
- Feynhiggs has more advanced one-loop corrections
- CPH (CPV version of SUBHPOLE) is more precise at the two-loop level

In each parameter space point the most conservative result is used

All implemented in HZHA with ISR and interference between identical final states from Higgstrahlung and boson fusion process

CPX scan

LHWG-Note/2004-01



● No lower $m_{H_{1,2}}$ limit

m_t dependent 95%CL on $\tan\beta$:

$m_t(\text{GeV}/c^2)$	Expected $\tan\beta$	Observed $\tan\beta$
174.3	>2.9	>3.0
179.3	>2.6	>2.7
183.0	>2.5	>2.5

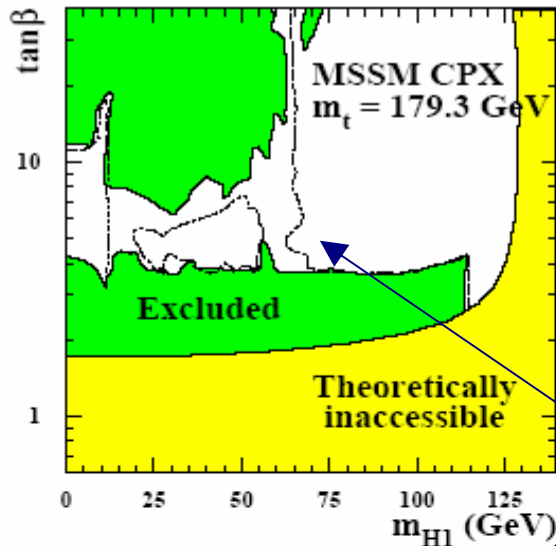
$m_t = 178.0 \pm 4.3 \text{ GeV}$ CDF & D0

For heavy H_2 :

- $H_1 \approx H_{SM}$
- $m_{H_1} \gtrsim 115 \text{ GeV}$

CPX scan

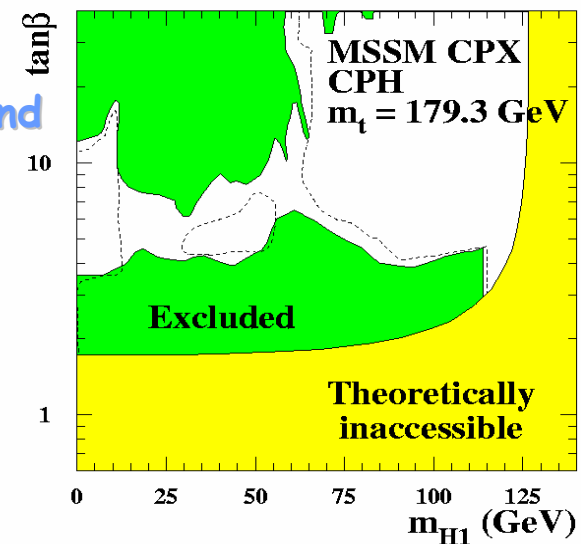
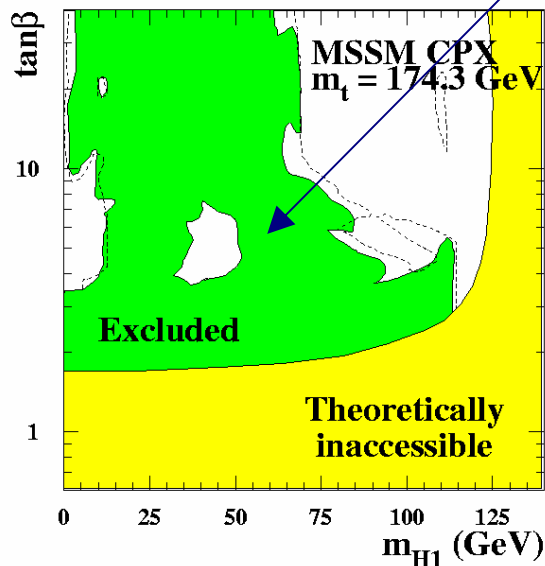
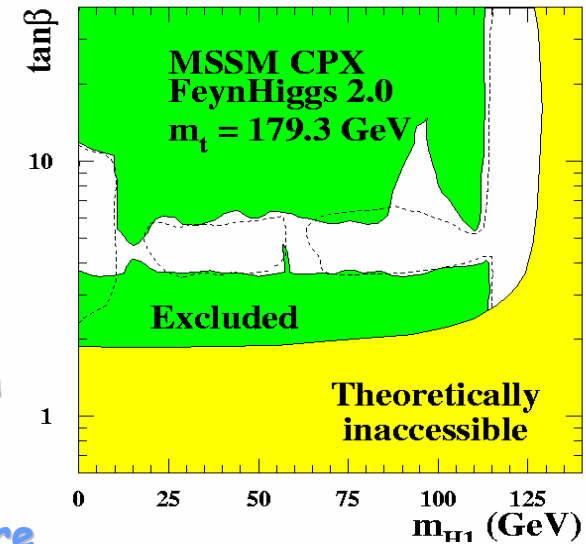
LHWG-Note/2004-01



Strong reduction of exclusion increasing $m_t \Rightarrow$ CP-violation:

for example $4 < \tan\beta < 10$, where both H_1Z & H_2Z are open

No excess larger than 3σ found



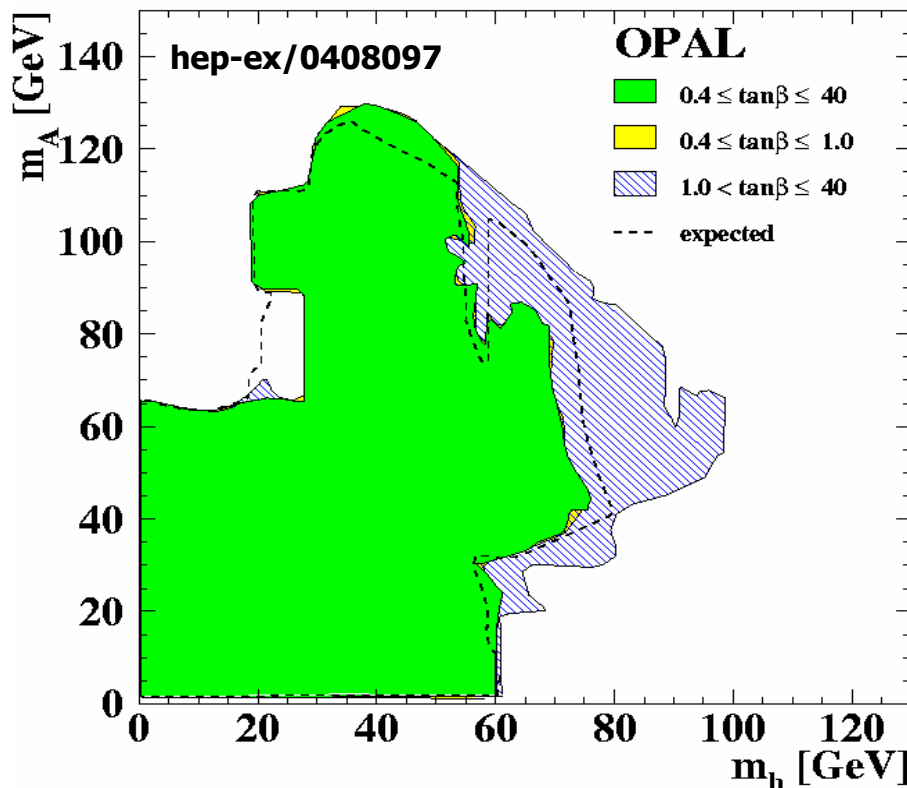
OPAL general 2HDM(II) scan

2HDM(II) have no constraint derived from SUSY:

- $h \rightarrow bb$ is not dominant decay, e.g for $\alpha=0$ $BR(h \rightarrow bb/\tau\tau)=0 \Rightarrow$ flavour-indep. searches
- large regions of the parameter space cannot be excluded by LEP



light Higgs not ruled out



Signal generated with HZHA

Free parameters:

- $1 < m_h < 130$ GeV
- 3 GeV $< m_A < 2$ TeV
- $\alpha = \pm \pi/2, \pm \pi/4, 0$
- $0.4 < \tan\beta < 40$

m_H and m_{H^\pm} kinematically inaccessible

Excluded rectangular region for $1 < m_h < 55$ GeV when $3 < m_A < 63$ GeV
But for $m_A > 63$ values of m_h down to 0 are still allowed !

No $\tan\beta$ exclusion independent of m_h/m_A

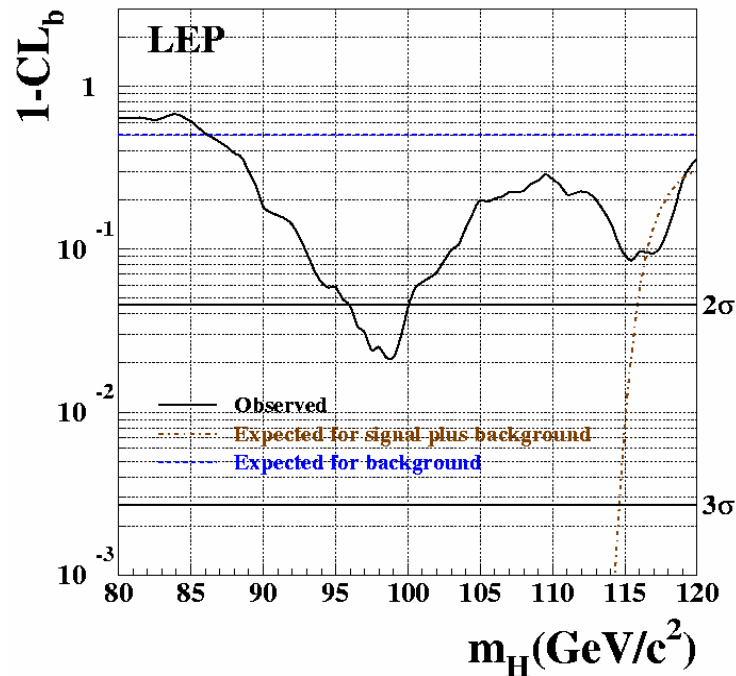
Conclusions

- At LEP we have searched for Higgses in several extensions of SM:
 - CP-conserving MSSM
 - CP-violating MSSM
 - 2HDM
- No evidence of the presence of a signal has been found
- Still there is room for the presence of a light Higgs:
 - In CP-violating MSSM
 - In 2HDMs
- Some theoretical papers interpreting small data-background discrepancy (about 2σ) in the context of specific CPC MSSM scenarios (require quite some tuning) as well as CPV MSSM and 2HDM.



Back-up

SM vs MSSM excess



The excesses at $m_h=98$ & 115 GeV Observed in MSSM are the same that where observed in the SM searches:

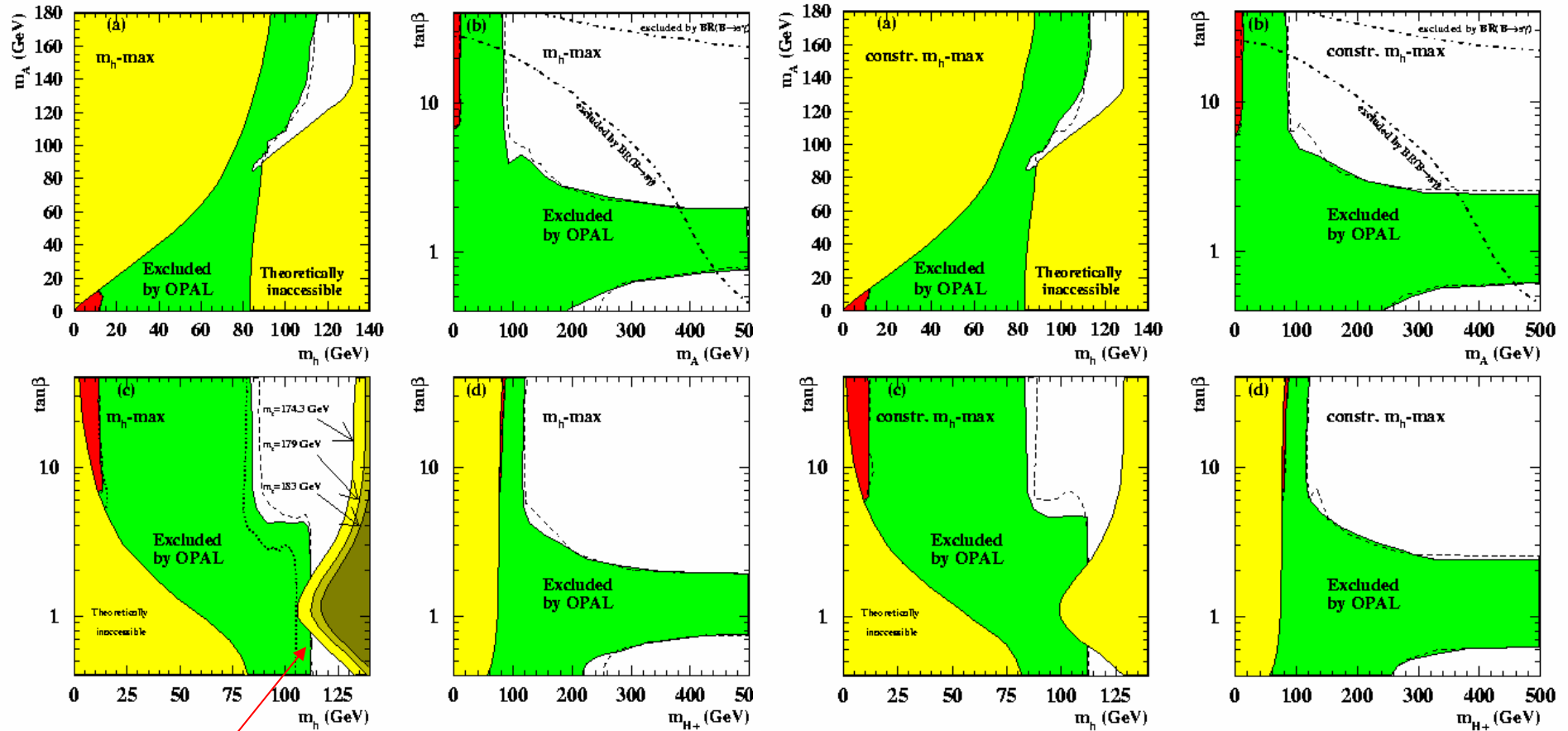
■ $m_h \sim 98$ GeV 2.3σ excess $1-CL_b = 2\%$ from all hZ channels
Not Compatible with a SM signal.

■ $m_h \sim 115$ GeV 1.7σ excess $1-CL_b = 9\%$ from ALEPH hZ 4-jet channel
Compatible with a SM signal.

$1-CL_b$ gives the Probability of a local fluctuation of background.
Probability that a fluctuation appears anywhere within a certain mass range is given by:

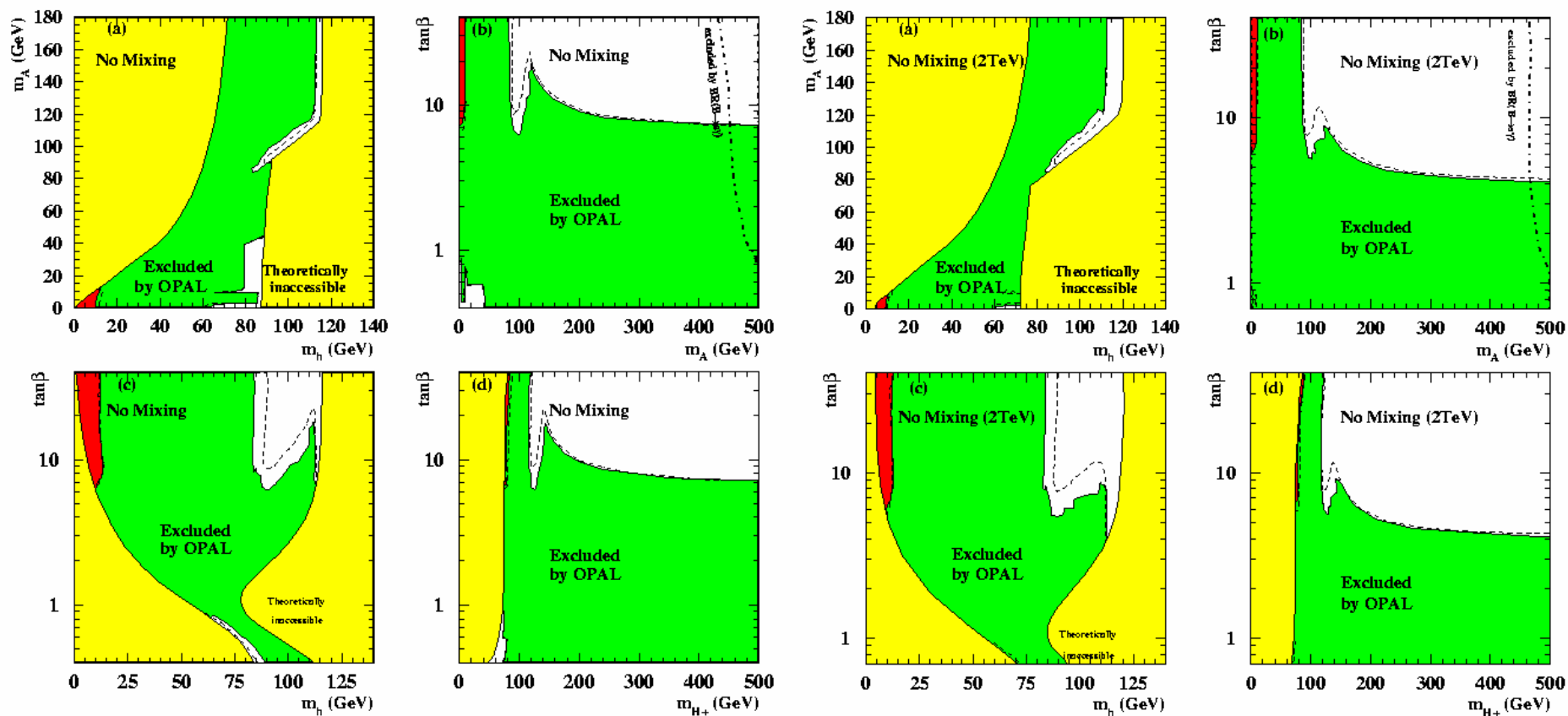
$$(1-CL_b) \times \frac{\Delta (\text{mass range})}{\text{mass resolution}}$$

m_h max scans

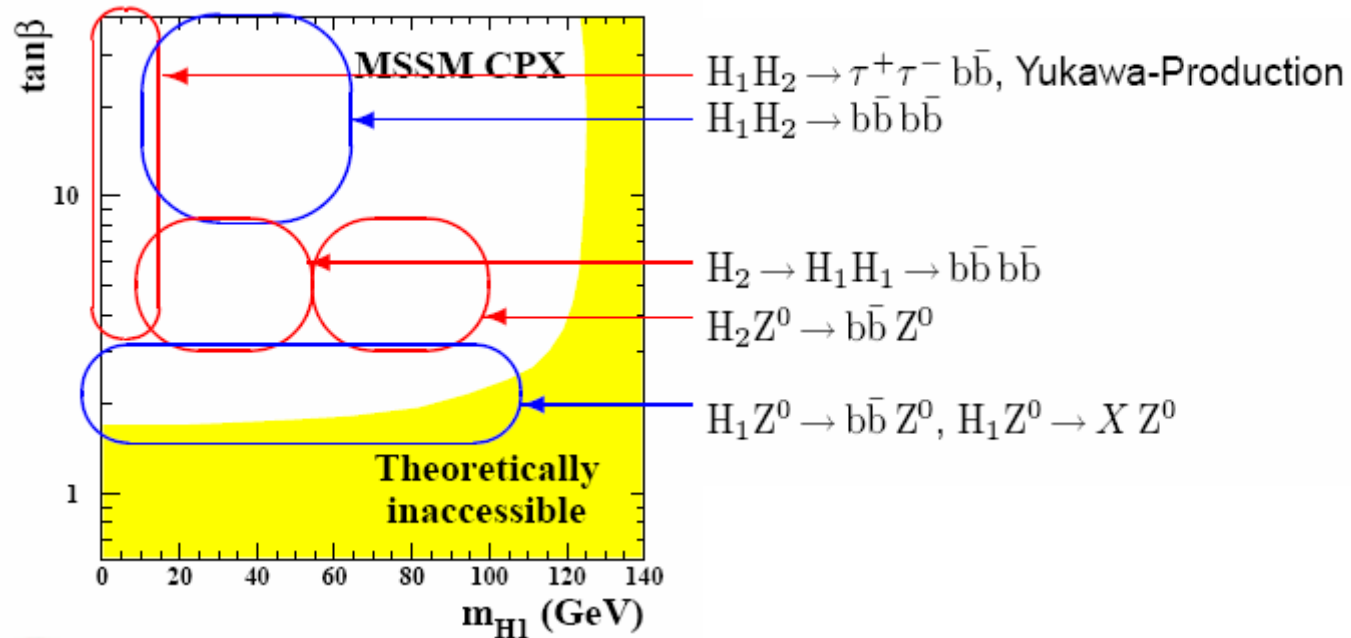


$m_h = 174.3, 179, 183$ GeV

No-mixing

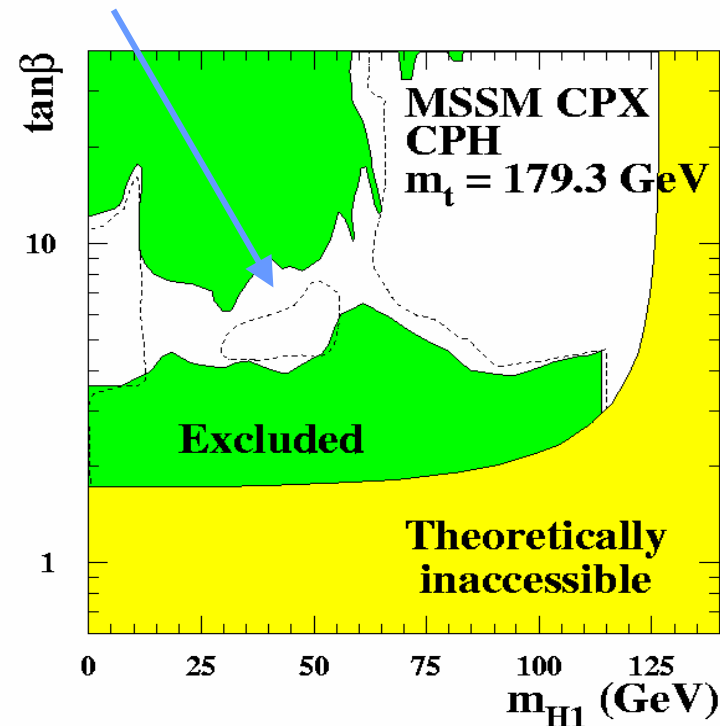
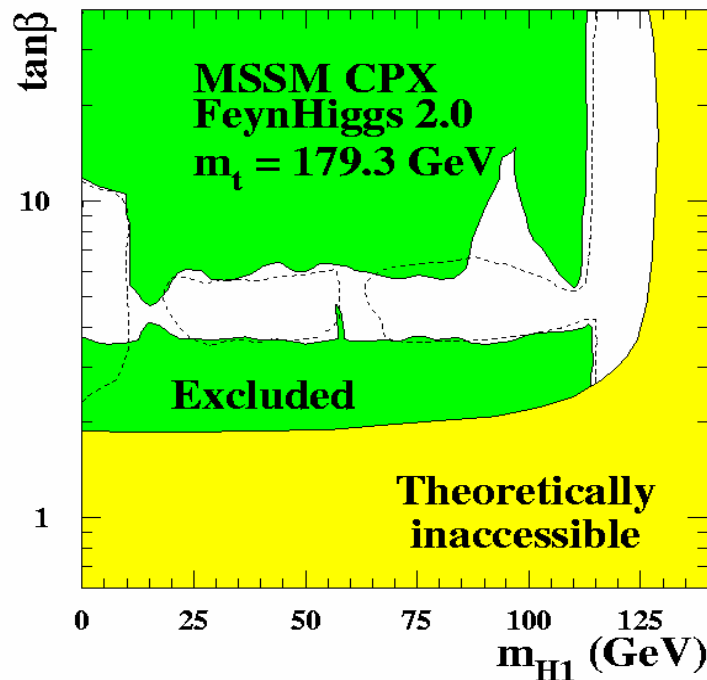


Useful searches for CPV MSSM

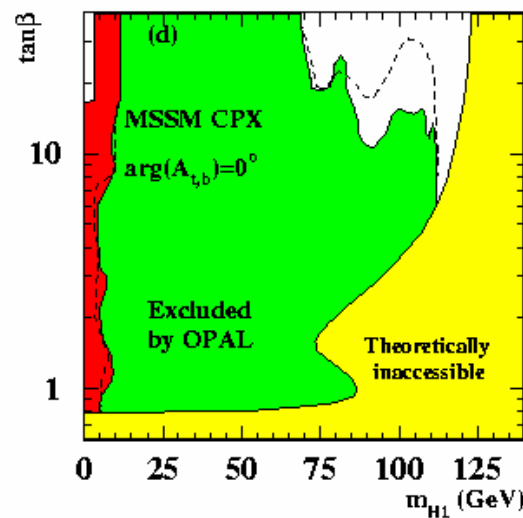
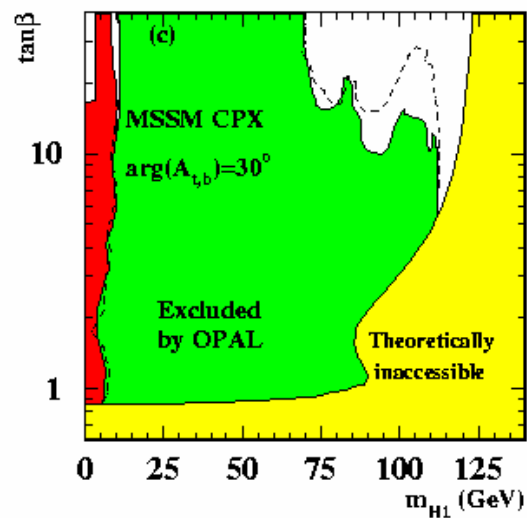
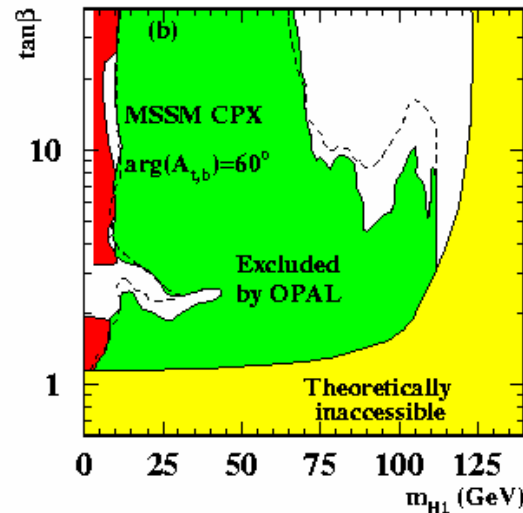
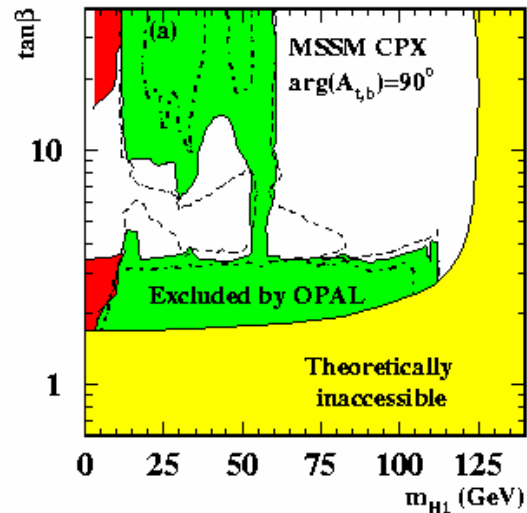


CPH vs FEYNHIGGS

- The largest discrepancy occurs for large $\tan\beta$ where Feynhiggs predicts a higher x-section for Higgstrahlung
- Data/background discrepancy in intermediate $\tan\beta$ region:
due to excess at $m_h \sim 98$ GeV which is the m_{H_2} mass in this region



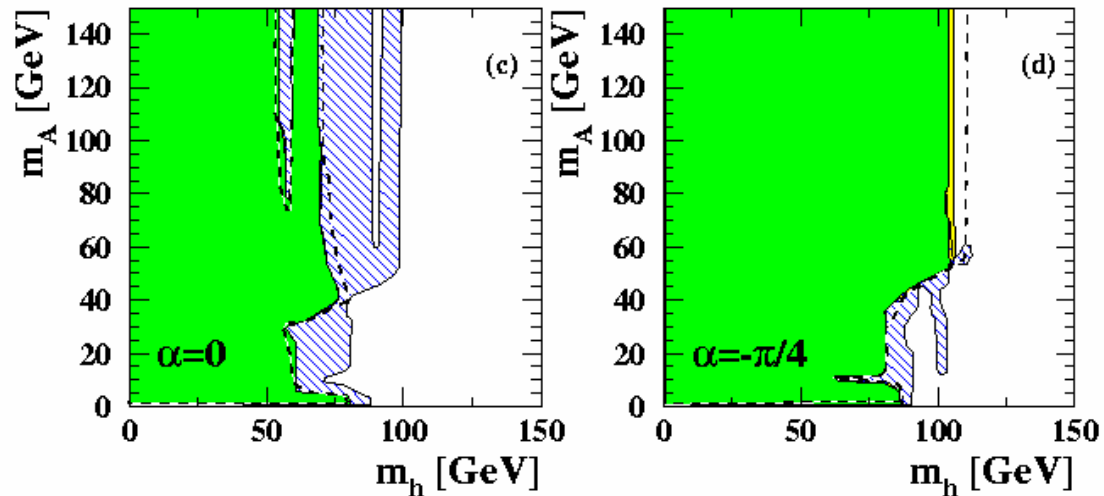
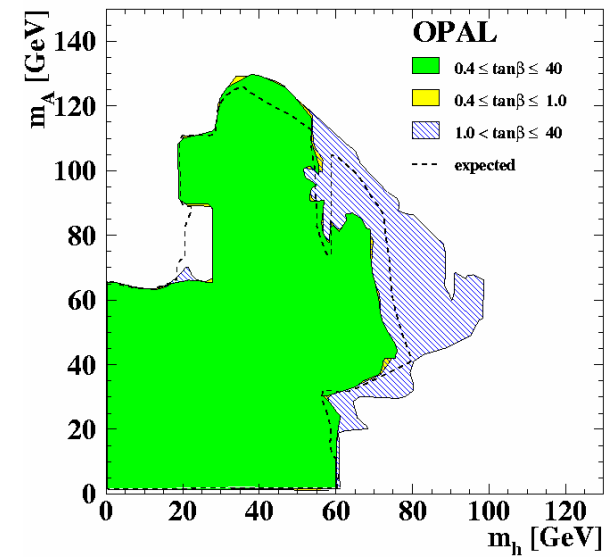
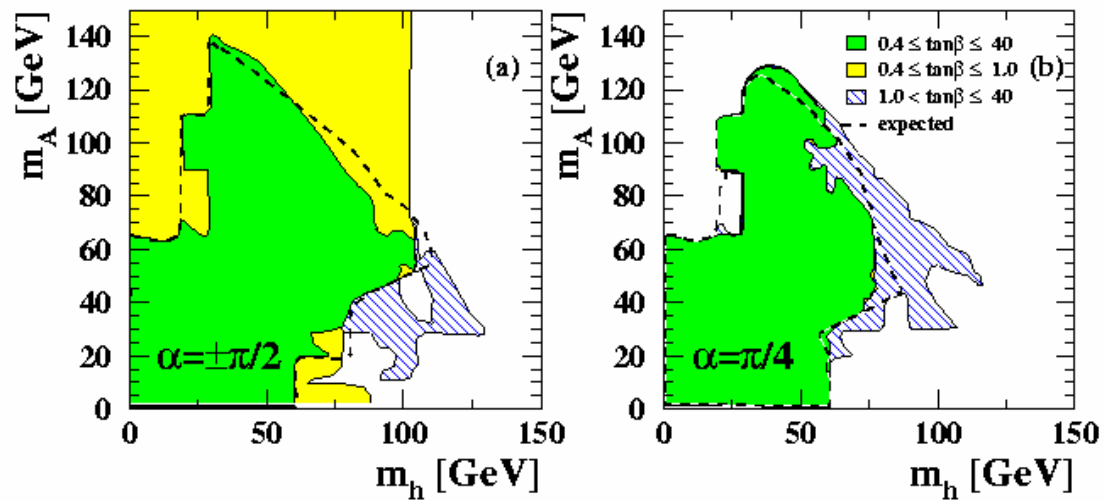
CPX-phases



Exclusion decreases
With increasing
 $\arg A_{t,b}$ (CP-Violation)

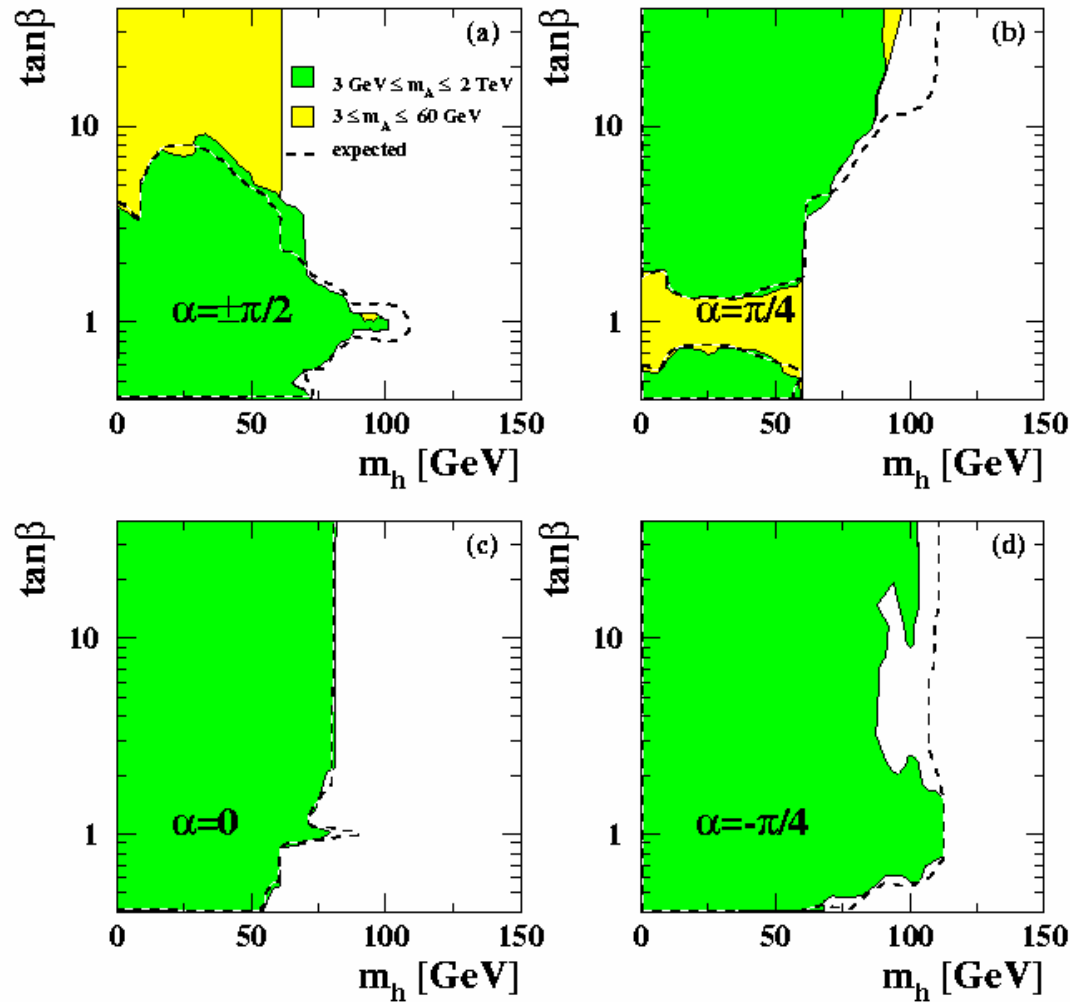
2HDM scan

OPAL



2HDM scan $\tan\beta$ exclusion

OPAL



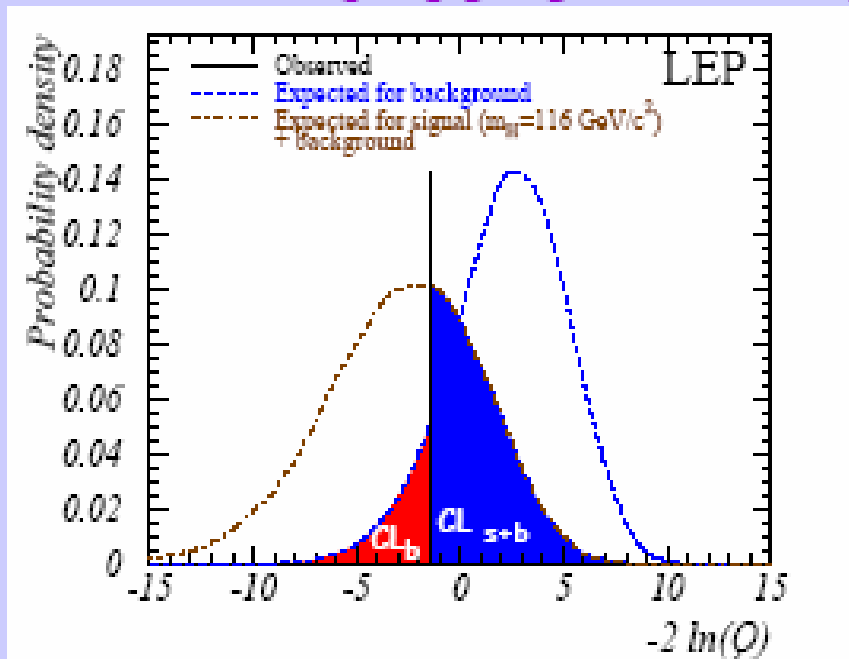
No $\tan\beta$ exclusion independent from m_h and m_A

The statistical method: "likelihood ratio"

- Data are subject to likelihood test ratio of 2 hypothesis
 - Data receive contributions from background processes only
 - Data receive contribution from background+SM Higgs boson signal of test mass m_H
- Channels are binned in 2-dimensional space (m_H^{rec}, \mathcal{G})

$$Q = \mathcal{L}_{s+b} / \mathcal{L}_b \quad \text{or} \quad -2 \ln Q = 2s_{tot} - 2 \sum_i n_i \ln(1 + \overset{\text{signal}}{s_i} / \underset{\text{background}}{b_i})$$

candidates
background



$1 - CL_{s+b}$ = measure of compatibility with s+b

$1 - CL_b$ = measure of incompatibility with b

$1 - CL_b$ 0.32 0.064 2.7×10^{-3} 6.3×10^{-5} 5.7×10^{-7}

1σ 2σ 3σ 4σ 5σ

$$CL_s = CL_{s+b} / CL_b$$