

## $K^+$ , $K^-$ and $K_S^0$ total rates in pp collisions at 24 GeV/c

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The knowledge of charged-kaon production rates in pNi or at least pp collisions at the beam momentum of 24 GeV/c is essential for proper tuning of the Monte Carlo generators for the DIRAC experiment. In particular it is important for estimation of the finite-size corrections to the lifetime of  $K^\pm\pi^\mp$ -atoms.

The available phase space coverage of the existing data in the center of mass energy range around  $\sqrt{s} = 6.8$  GeV is known to be incomplete and partially incompatible especially for charged kaons. The double differential cross sections for charged kaons in pp collisions at 24 GeV/c beam momentum have been measured by the CERN-Rome group, Allaby et al. [1, 2]. However, the precise data on the production of charged kaons in pp collisions at 158 GeV/c beam momentum recently obtained by the NA49 experiment at CERN [3] and a new evaluation of the energy dependence of kaon production in [3] has revealed that the resulting  $s$ -dependence for the  $K^+$  production looks unphysical indicating an excess of the order of 60% in the  $K^+$  yields of Allaby et al. For example, the Allaby et al. data at  $p_T = 0.632$  GeV/c, Feynman- $x_F=0.2$  (Fig. 66a in [3]) and  $x_F=0.2$ ,  $p_T = 0.55$  GeV/c (Fig. 66b in [3]) are on the same level as the NA49 data for 24 GeV/c beam momentum and even *higher* for 19.2 GeV/c.

Meanwhile the Fritiof 6 - Monte Carlo generator for the DIRAC experiment was tuned to the measurements [4] of particle production in proton interactions with nuclei at 24 GeV/c. This experiment [4] was performed with the same spectrometer of the CERN-Rome group as the one in the experiment of Allaby et al., with the hydrogen target just replaced by a set of nuclear targets. Therefore if an excess in the  $K^+$  yields of Allaby et al. is related to the performance of their spectrometer, the data of Eichten et al. [4] have to be also treated with a caution.

The  $K^+$  and  $K^-$  total rates in pp collisions at 24 GeV/c were determined in the bubble chamber experiment [5]. For this a sample of events with inclusive production of two strange particles with opposite strangeness was analyzed. Only those channels in which a charged kaon was produced in association with a neutral kaon, a lambda or a charged sigma, i.e. the final states  $K^+K_S^0X$ ,  $K_S^0K^-X$ ,  $K^+\Lambda X$ ,  $K^+\Sigma^+X$  and  $K^+\Sigma^-X$ , were considered. Charged kaons were detected either by their pionic or muonic decay modes or, in the most cases, by a *statistical procedure*. The significant pion contamination (of about 40% for  $p_{tab} > 1$  GeV/c) due to this separation procedure could indeed be not very important for investigation of net strangeness and strangeness-transfer distributions which represented the main task of paper [5]. However the systematic errors on the total  $K^+$  and  $K^-$  rates obtained in [5] and used in [3] can be quite large (as it is evident from Figs. 1 and 2 to be presented below).

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In this note, an attempt is made to estimate the  $K^+$  and  $K^-$  total rates in pp interactions at 24 GeV/c. It is based on a study of energy evolution of the total yields of the charged and neutral kaons in [3] and reliable value of the measured total  $K_S^0$  rate in the bubble chamber pp experiment at 24 GeV/c [6].

Investigation of the energy dependence of the total charged kaon rates in [3] has been based on assumption that this dependence is closely related to the one for the neutral kaons:

$$\langle K^+ \rangle + \langle K^- \rangle = \langle K^0 \rangle + \langle \bar{K}^0 \rangle = 2\langle K_S^0 \rangle$$

or

$$R = 0.5(\langle K^+ \rangle + \langle K^- \rangle) / \langle K_S^0 \rangle = 1, \quad (1)$$

as might be expected from isospin symmetry. This relation is not fulfilled for  $\phi$  with its branching fractions  $Br(\phi \rightarrow K\bar{K}) = 0.491$  and  $Br(\phi \rightarrow K_S^0 K_L^0) = 0.341$  [7]. However the  $\phi$  total production rate in pp collisions at 24 GeV/c is relatively small  $\langle \phi \rangle = 0.0052 \pm 0.0011$  [8]. Therefore the influence of the  $\phi$  production on (1) is negligible.

Indeed, the ratio R was determined in [3] to be 2.8 at  $\sqrt{s} = 3$  GeV, 1.4 at  $\sqrt{s} = 3.5$  GeV and 1.27 at  $\sqrt{s} = 4$  GeV. Thus it approaches unity rather quickly with increasing energy, so that it may be assumed, within a few percent error margin, that  $R = 1$  at  $\sqrt{s} > 5$  GeV.

In Fig. 1 taken from [3] (Fig. 130), the total yields  $\langle K^+ \rangle$ ,  $\langle K^- \rangle$  and  $\langle K_S^0 \rangle$  are shown as a function of  $\sqrt{s}$ . The full line through the  $K_S^0$  data points is an eyeball fit which gives a consistent description of the data within point-by-point deviation of typically 10-20%. In a second step the ratios  $\langle K^+ \rangle / \langle K_S^0 \rangle$  and  $\langle K^- \rangle / \langle K_S^0 \rangle$  were obtained in [3] from the available data. These ratios are presented in Fig. 2 also taken from [3] (Fig. 131). A smooth  $\sqrt{s}$  dependence was imposed on these data points between the low energy range  $\sqrt{s} < 4.8$  GeV and the higher energies  $\sqrt{s} > 6.8$  GeV since the cross sections in the 4.8-6.8 GeV range were shown in [3] to deviate upwards. The  $\langle K^+ \rangle / \langle K_S^0 \rangle$  and  $\langle K^- \rangle / \langle K_S^0 \rangle$  ratios thus obtained were then used to produce smooth lines in Fig. 1 through the  $K^+$  and  $K^-$  data by multiplying with the  $K_S^0$  interpolation.

From the smooth lines in Fig. 2 one has at  $\sqrt{s} = 6.8$  GeV

$$\langle K^+ \rangle / \langle K_S^0 \rangle = 1.51 \quad \text{and} \quad \langle K^- \rangle / \langle K_S^0 \rangle = 0.43. \quad (2)$$

With the  $K_S^0$  total rate determined in the bubble chamber pp experiment at 24 GeV/c beam momentum [6]

$$\langle K_S^0 \rangle = 0.0507 \pm .0020^2 \quad (3)$$

one then obtains:

$$\langle K^+ \rangle = 0.0766 \quad \text{and} \quad \langle K^- \rangle = 0.0218 \quad (4)$$

with

$$\langle K^+ \rangle / \langle K^- \rangle = 3.52. \quad (5)$$

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<sup>2</sup>Different data points at  $\sqrt{s} = 6.8$  GeV in Figs. 1 and 2 represent, in fact, the results of the same experiment. The data points with the lowest  $\langle K_S^0 \rangle$  in Fig. 1 and, respectively, the highest  $\langle K^+ \rangle / \langle K_S^0 \rangle$  and  $\langle K^- \rangle / \langle K_S^0 \rangle$  in Fig. 2 have to be ignored since they are obtained from earlier result [9] of this experiment on the  $K_S^0$  rate. The final result was given in [5, 6].

The errors of the values (4) and (5) can be roughly estimated as equal to 10-20%. It is therefore quite assuring that the ratio of the  $K^{*+}(892)$  and  $K^{*-}(892)$  total rates in the same pp experiment at 24 GeV/c [6] is exactly the same as the  $K^+/K^-$  ratio:

$$\langle K^{*+}(892) \rangle / \langle K^{*-}(892) \rangle = 3.5 \pm 0.8, \quad (6)$$

as could be expected.

Significant difference between  $K^+$  and  $K^-$  production arising from associate kaon plus hyperon versus kaon pair production persists even at SPS energies [3]. Therefore it is important to study evolution of the  $K^+/K^-$  ratio as a function of the kinematic variables and its  $s$ -dependence.

Besides, possible signature of new physics by the creation of a deconfined state of matter in heavy nuclear interactions, looked for since a long time, is supposed to rely completely on a comparison with elementary or proton-nuclei collisions. Therefore the detailed study of behaviour of the kaon production from low energy up to RICH and collider energies in pp and proton-nuclei collisions appears to be very important.

From this point of view it would be interesting to obtain the precise data on the  $K^+/K^-$ ,  $\bar{p}/p$  and  $\bar{\Lambda}/\Lambda$  ratios and their dependence on Feynman- $x_F$  and transverse momentum variables in pNi collisions at 24 GeV/c as it seems possible from the data already collected by the DIRAC experiment.

## References

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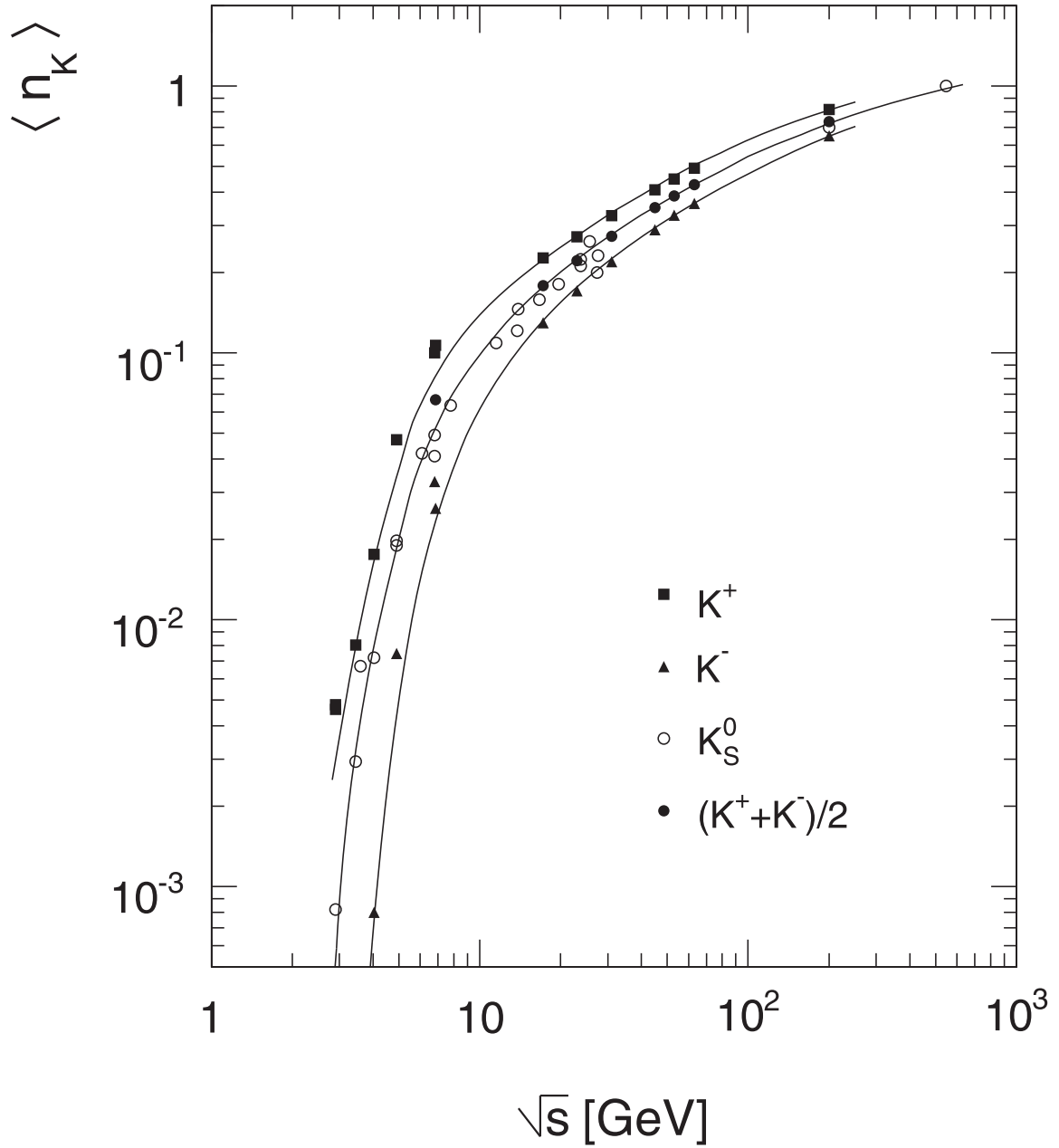


Figure 1: Fig. 130 taken from [3]: Total yields  $\langle K^+ \rangle$ ,  $\langle K^- \rangle$  and  $\langle K_S^0 \rangle$  as a function of  $\sqrt{s}$ . The full line through the  $K_S^0$  results is eyeball fit, the lines through the  $K^+$  and  $K^-$  data are derived from Fig. 2. The full circles in the  $K_S^0$  data correspond to  $0.5(\langle K^+ \rangle + \langle K^- \rangle)$  established at corresponding  $\sqrt{s}$  values

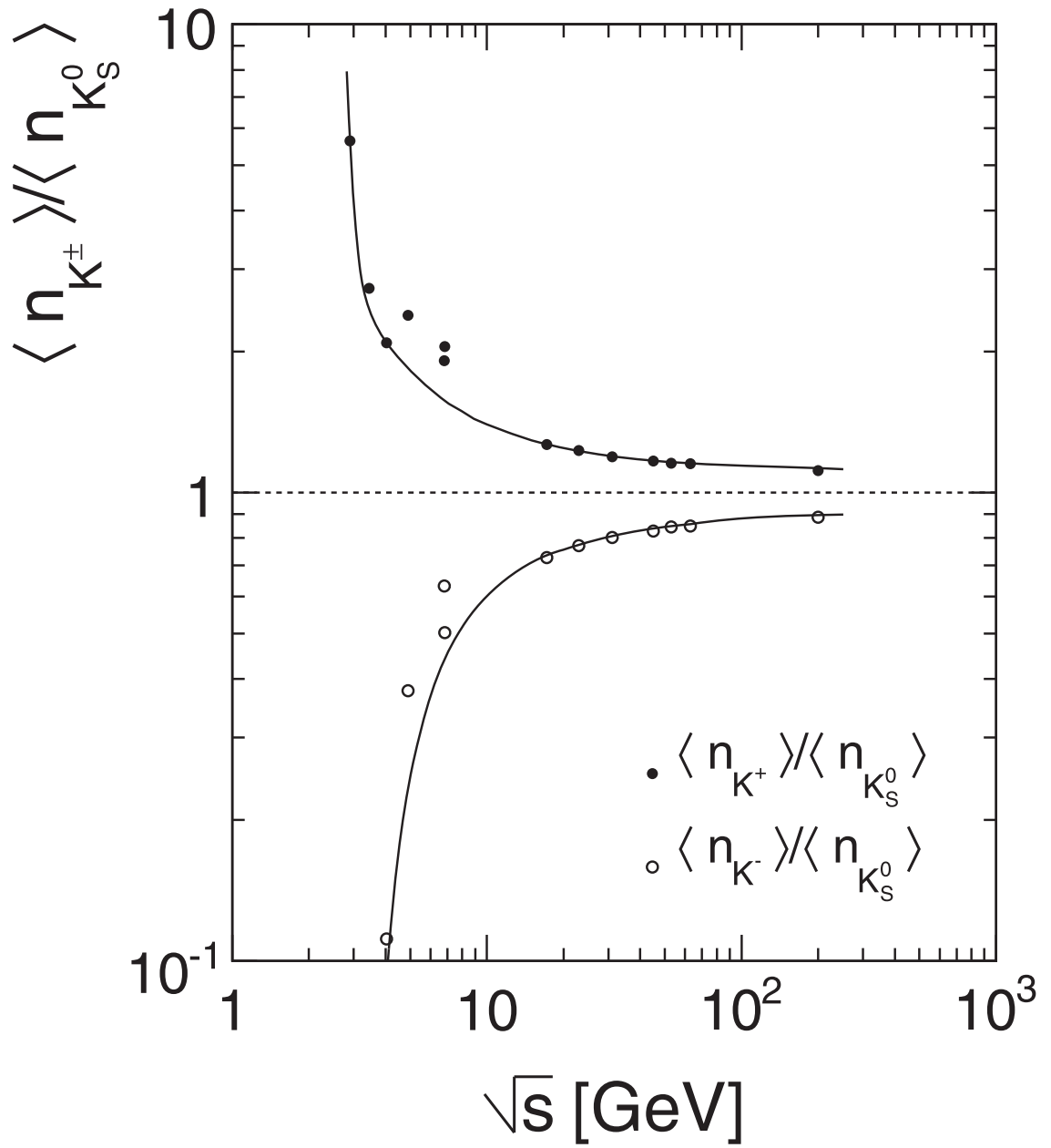


Figure 2: Fig. 131 from [3]: Ratios  $\langle K^+ \rangle / \langle K_S^0 \rangle$  and  $\langle K^- \rangle / \langle K_S^0 \rangle$  as a function of  $\sqrt{s}$ . The full lines are eyeball fits through the data at  $\sqrt{s} < 4.9$  GeV and  $\sqrt{s} > 6.8$  GeV