

Calculation of \mathcal{V}_{fin} in the FKS framework: an example using POWHEG-BOX scales

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Abstract

We show how to calculate the \mathcal{V}_{fin} term needed by the POWHEG-BOX program, starting from a generic expression for the virtual term \mathcal{V}_b . We show this in the case of DY processes.

We start from FNO, eq.7.190, which is taken from the original paper by Altarelli, Ellis and Martinelli:

$$\mathcal{V}_{b, q\bar{q}} = \left(\frac{4\pi\mu_{\text{R}}^2}{s} \right)^\epsilon \frac{\Gamma(1+\epsilon)\Gamma(1-\epsilon)^2}{\Gamma(1-2\epsilon)} \frac{\alpha_{\text{S}}}{2\pi} C_{\text{F}} \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} - 8 + \pi^2 \right] \mathcal{B}_{q\bar{q}}, \quad (1)$$

Here μ_{R} stands for the renormalization scale, and $s = 2(k_{\oplus} \cdot k_{\ominus})$. To extract \mathcal{V}_{fin} from eq. (1), we need to equal the expression in (1) to the following one (FNO, eq.2.92):

$$\mathcal{V}_b = \mathcal{N} \frac{\alpha_{\text{S}}}{2\pi} \left[- \sum_{i \in \mathcal{I}} \left(\frac{1}{\epsilon^2} C_{f_i} + \frac{1}{\epsilon} \gamma_{f_i} \right) \mathcal{B} + \frac{1}{\epsilon} \sum_{\substack{i, j \in \mathcal{I} \\ i \neq j}} \log \frac{2k_i \cdot k_j}{Q^2} \mathcal{B}_{ij} + \mathcal{V}_{\text{fin}} \right] \quad (2)$$

where

$$\mathcal{N} = \frac{(4\pi)^\epsilon}{\Gamma(1-\epsilon)} \left(\frac{\mu_{\text{R}}^2}{Q^2} \right)^\epsilon. \quad (3)$$

For the case at hand, equation (2) reduces to

$$\mathcal{V}_{b, q\bar{q}} = \left(\frac{4\pi\mu_{\text{R}}^2}{Q^2} \right)^\epsilon \frac{1}{\Gamma(1-\epsilon)} \frac{\alpha_{\text{S}}}{2\pi} \left[- \left(\frac{2}{\epsilon^2} + \frac{3}{\epsilon} - \frac{2}{\epsilon} \log \frac{s}{Q^2} \right) C_{\text{F}} \mathcal{B}_{q\bar{q}} + \mathcal{V}_{\text{fin}, q\bar{q}} \right]. \quad (4)$$

where we used the fact that the only colored partons involved are the two incoming quarks, and we also have that $\mathcal{B}_{\oplus\ominus}^{q\bar{q}} = C_{\text{F}} \mathcal{B}_{q\bar{q}}$.

We notice that Q^2 is here an *arbitrary* scale. In the FKS approach, the soft-virtual term is obtained summing to \mathcal{V}_{fin} a quantity that is Q^2 dependent (see FNO, eq. 2.99, 2.100 and 2.101). In the POWHEG-BOX code, this operation is automatized and it is performed in a file that the user is not supposed to edit. Therefore, an explicit choice for this scale was needed. Since we choosed

$$Q^2 = \mu_{\text{R}}^2, \quad (5)$$

it is important to use the same value for Q^2 in eq. (4), in order to extract a consistent expression for \mathcal{V}_{fin} .

Therefore we evaluate the expression in (4) with $Q^2 = \mu_{\text{R}}^2$, and equal it to the r.h.s. of eq. (1), obtaining:

$$(4\pi)^\epsilon \frac{1}{\Gamma(1-\epsilon)} \left[- \left(\frac{2}{\epsilon^2} + \frac{3}{\epsilon} - \frac{2}{\epsilon} \log \frac{s}{\mu_{\text{R}}^2} \right) \mathcal{B}_{q\bar{q}} + \frac{\mathcal{V}_{\text{fin}, q\bar{q}}}{C_{\text{F}}} \right] = \left(\frac{4\pi\mu_{\text{R}}^2}{s} \right)^\epsilon \frac{\Gamma(1+\epsilon)\Gamma(1-\epsilon)^2}{\Gamma(1-2\epsilon)} \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} - 8 + \pi^2 \right] \mathcal{B}_{q\bar{q}}. \quad (6)$$

By using the fact that

$$\frac{\Gamma(1+\epsilon)\Gamma(1-\epsilon)^2}{\Gamma(1-2\epsilon)} = \frac{1}{\Gamma(1-\epsilon)} + \mathcal{O}(\epsilon^3), \quad (7)$$

we have

$$\frac{\mathcal{V}_{\text{fin}, q\bar{q}}}{C_{\text{F}}} = \left[\left(\frac{\mu_{\text{R}}^2}{s} \right)^\epsilon \left(-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} - 8 + \pi^2 \right) + \left(\frac{2}{\epsilon^2} + \frac{3}{\epsilon} - \frac{2}{\epsilon} \log \frac{s}{\mu_{\text{R}}^2} \right) \right] \mathcal{B}_{q\bar{q}}, \quad (8)$$

that reduces to

$$\mathcal{V}_{\text{fin}, q\bar{q}} = \mathcal{B}_{q\bar{q}} \left[\pi^2 - 8 - 3 \log \frac{\mu_{\text{R}}^2}{s} - \log^2 \frac{\mu_{\text{R}}^2}{s} \right] C_{\text{F}} \quad (9)$$

after expanding $(\mu_{\text{R}}^2/s)^\epsilon$ as usual. The expression in eq. (9) is the one needed by the POWHEG-BOX program, and it has to be coded inside the subroutine `setvirtual`.