

On the predictivity of single field inflationary models.

Subodh P. Patil

CERN

HKUST IAS, May 29th 2014

Based on C.P. Burgess, S.P. Patil, M. Trott, arXiv:1402.1476; to appear, JHEP

What we know:

The word is in! (ACT, Planck, SPT) Spectacular confirmation of the (six parameter) phenomenological Λ CDM model.

- ▶ Assuming $\Omega_{tot} = 1, w_\Lambda = -1, \sum_i m_\nu = 0 \dots$

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- ▶ $\Omega_c h^2 = 0.1196 \pm 0.0031 \quad \ln(10^{10} A_s) = 3.103 \pm 0.072$
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- ▶ \exists a *single* effectively light degree of freedom at $\sim \epsilon^{1/4} 10^{16} \text{ GeV}$. BICEP2 $\rightarrow 10^{16} \text{ GeV}$?
 - ▶ whose field modes began in the relevant vacuum state (BD)
 - ▶ whose self interactions and interactions with other fields are sufficiently weak or irrelevant *throughout* inflation
 - ▶ which at the same time couples strongly enough to some sector that contains the standard model so that efficient (pre)heating occurs...

Effective field theory

How do these requirements sit with our understanding of Effective field theory? Obtaining *enough* inflation requires field excursions.

▶ $\mathcal{L} = -\frac{1}{2} (\partial\phi)^2 - V(\phi) + \dots$, where $\dots = \sum_i c_i \frac{\mathcal{O}_i(\phi)}{\Lambda_i^{D_i-4}}$

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- ▶ By thinking about this problem honestly, can rule out a lot of models a priori.

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- ▶ Not predictive \leftrightarrow need to know UV completion.

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- ▶ N.B. kinetic mixing of the radial mode (singlet) and the Goldstone modes! – crucially distinction between Higgs inflation and singlet scalar field w/ quartic potential.

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In unitary gauge, $H = (0, \frac{h}{\sqrt{2}})^T$ – we define $\langle h \rangle = \frac{\bar{\phi}}{\sqrt{2}}$, with $\bar{\phi}$ denoting classical the background field expectation value for the Higgs.

- ▶ ... satisfies the boundary condition $\bar{\phi}_{ew} = v \equiv 246 \text{ GeV}$

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$$V_E(\chi) = \frac{\lambda M_{pl}^4}{4\xi^2} \left[1 - e^{\frac{-2\chi}{\sqrt{6}M_{pl}}} \right]^2 + \dots$$

▶ Exponentially flat in the large field region $\bar{\phi} \gg M_{pl}/\xi$, slow roll conditions satisfied at $\bar{\phi} \gtrsim M_{pl}/\sqrt{\xi}$.

▶ ... denotes corrections to the classical tree level potential.

▶

$$n_s = 1 - 6\epsilon + 2\eta, \quad r = 16\epsilon; \quad \epsilon(\bar{\phi}) \simeq \frac{4M_{pl}^4}{3\bar{\phi}^4\xi^2}, \quad \eta(\bar{\phi}) \simeq \frac{4M_{pl}^4}{3\bar{\phi}^4\xi^2} \left(1 - \xi \frac{\bar{\phi}^2}{M_{pl}^2} \right)$$

Higgs Inflation

In unitary gauge, $H = (0, \frac{h}{\sqrt{2}})^T$ – we define $\langle h \rangle = \frac{\bar{\phi}}{\sqrt{2}}$, with $\bar{\phi}$ denoting classical the background field expectation value for the Higgs.

▶ ... satisfies the boundary condition $\bar{\phi}_{ew} = v \equiv 246 \text{ GeV}$

▶ Canonically normalized Higgs field $h \rightarrow \chi$, with

$$\frac{d\chi}{dh} = \frac{[1 + (\xi + 6\xi^2)(h/M_{pl})^2]^{1/2}}{1 + \xi(h/M_{pl})^2}.$$

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▶ For ~ 58 e-folds, $n_s \simeq 0.967$, $r \simeq 0.0031$

Unitarity bounds

An EFT is only defined up to the scale at which unitarity is violated, or the cut-off, whichever is the lower of the two.

- ▶ In the small field regime: $\Lambda_{ew} \simeq \frac{M_{pl}}{\xi}$, $v \leq \bar{\phi} \leq \frac{M_{pl}}{\xi}$

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Bezrukov et al 2010

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- ▶ You are *only just* OK! However, you will be riding just below the floating cut-off all throughout RG running.
- ▶ 'Threshold' effects could affect your observables if you want to connect to low energy EW physics (one of the major attractions of Higgs Inflation).

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To make predictions for CMB observables, we have to compute the effective potential in the inflationary regime.

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de Simone, Hertzberg, Wilczek 2008; Bezrukov, Mangin, Shaposhnikov 2008

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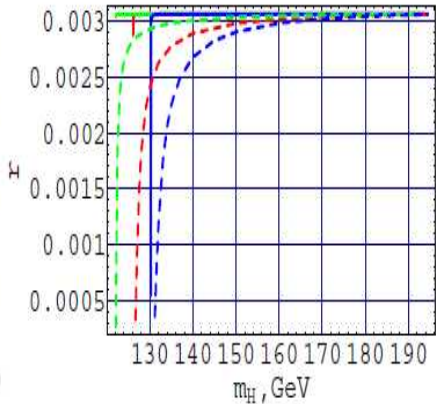
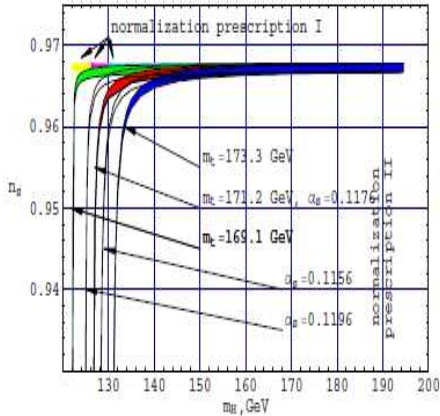
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- ▶ Discovery of the Higgs 2012 \rightarrow have to take SM parameters up to three sigma away from their central values to obtain positive quartic coupling at $M_{pl}/\sqrt{\xi}$.

RGE Improved potential



► From Bezrukov 2013.

Running up to inflation

On the predictivity of single field inflationary models.

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In order to implement the running honestly, had to run up from the SM parameters at top mass (computed at 2 loops to NNLO):

Buttazzo, Degrassi, Giardino, Giudice, Salvio, Strumia 2013

$$\blacktriangleright \lambda(m_t) = 0.12711, y_t(m_t) = 0.93558, g'(m_t) = 0.35761, \\ g(m_t) = 0.64822, g_s(m_t) = 1.1666, \xi_0 = 2300 + \delta\xi$$

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- ▶ But what about the fact that $\Lambda_{int} \simeq 4\pi\bar{\phi}$ during inflation?
Threshold effects at M_{pl}/ξ ?

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From m_t to E_{inf}

Higgs inflation is not renormalizable. Obligated to treat it as an effective description. D=6 SM operators mix into the running of SM parameters under RGE, even at one loop. [Jenkins, Manohar, Trott 2013](#)

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- ▶ Exact expressions requires computation of 59×59 D anomalous dimension matrix. [Jenkins, Manohar, Trott 2013](#)
- ▶ Wilson coefficients affect the running substantially. Unless we know the UV completion (i.e. can specify all the coefficients), have to allow for them to represent a *theoretical uncertainty* in the predictions of Higgs inflation.

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Can repeat the analysis done by others allowing for the theoretical uncertainty represented by these unknown Wilson coefficients, and compute the effects on CMB observables. [C.P.Burgess, S.P.Patil, M.Trott,](#)

[arXiv:1402.1476](#)

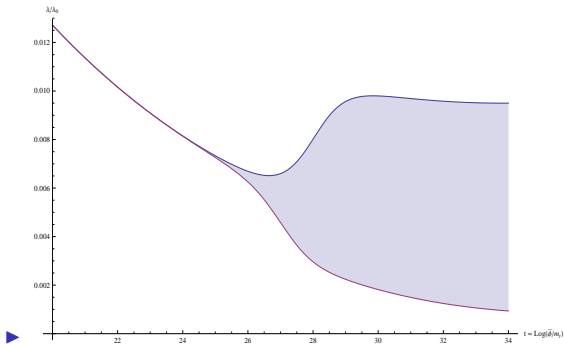


Figure : The effect of the unknown UV completion on the running of the quartic coupling in the Higgs inflation scenario.

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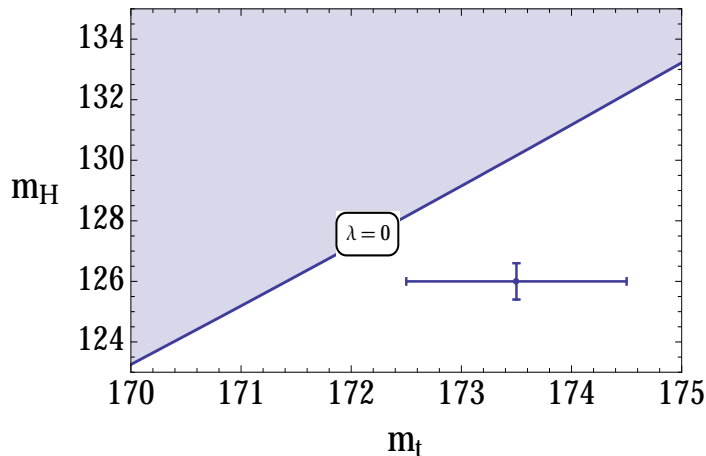


Figure : Initial conditions that separate $\lambda < 0$ from $\lambda > 0$ at $M_{pl}/\sqrt{\xi}$

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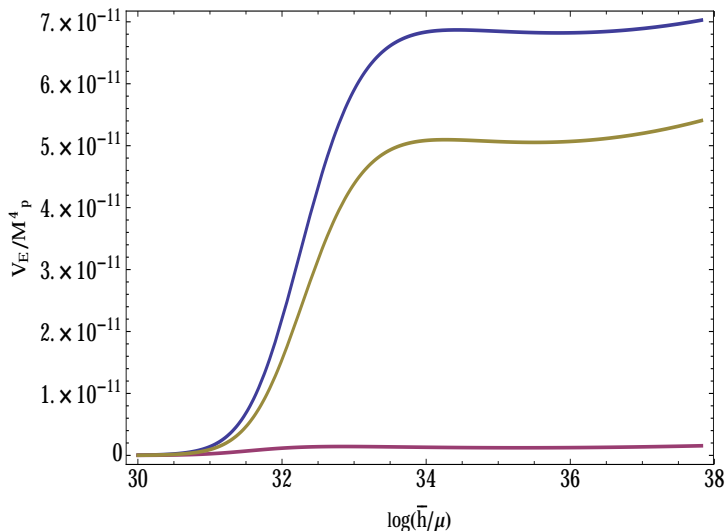


Figure : Effective potential over the range of Wilson coefficients

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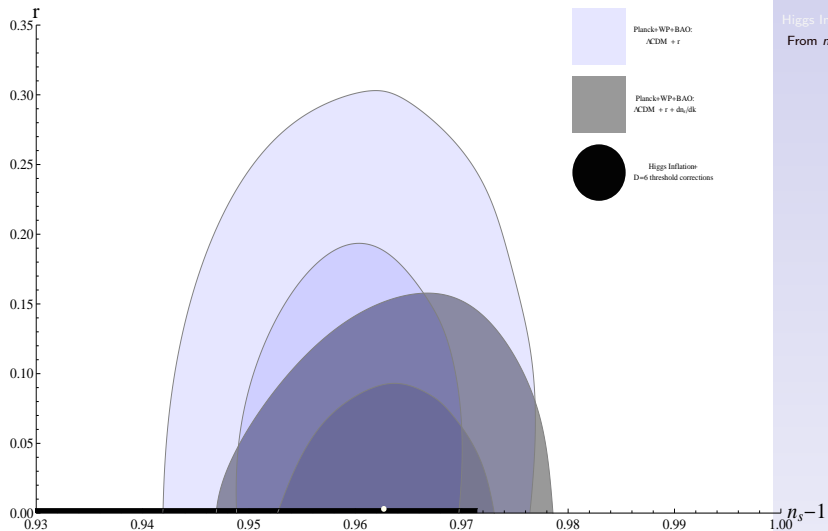
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- ▶ Classically can reproduce the outcome of almost any observations.

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- ▶ Are there any mechanisms to generate inflation that are not?
- ▶ A related question to predictivity– is inflation falsifiable?